

1 Waveform inversion for S-wave structure in the
2 lowermost mantle beneath the Arctic: Implications
3 for mineralogy and chemical composition

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4 We perform waveform inversion for the radial profile of shear wave veloc-
5 ity in the lowermost mantle beneath the Arctic. We use waveforms from the
6 CANOE (CANadian NOrthwest Experiment) array, which greatly enhances
7 the resolution in the lowermost mantle as compared to earlier studies. We
8 find a velocity increase in at depths from 2500 to 2700 km and a velocity de-
9 crease at depths from 2700 km to the core-mantle boundary (CMB). We in-
10 terpret the velocity increase as associated with the phase transition from per-
11 ovskite (pv) to post-perovskite (ppv), and the velocity decrease as due to a
12 temperature increase in the thermal boundary layer. The shear wave veloc-
13 ity immediately above the core-mantle boundary (CMB) is 7.11 km/s, while

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14 that beneath Central America is 7.25 km/s. This suggests that the propor-
15 tion of impurities in Mg-pv or Mg-ppv beneath the Arctic is 6 mol% larger
16 than that beneath Central America.

1. Introduction

17 Information on the elastic structure of the lowermost mantle can contribute to under-
18 standing geodynamics and the earth's evolution. It is generally expected that there is
19 complex structure in the lowermost mantle (just as in the uppermost mantle), because of
20 the thermal boundary layer at the base of the mantle. There is chemical fractionation in
21 the lowermost mantle due to effects such as the interaction between the mantle and the
22 core and partial melting. The most abundant material in the lower mantle is MgSiO_3 ,
23 which undergoes a phase transition from perovskite (pv) to post-perovskite (ppv) near
24 the core-mantle boundary (CMB) [Murakami *et al.*, 2004]. Seismological data can pro-
25 vide constraints on these phase transitions and also on the possible presence of impurities.

26 Many tomographic studies have revealed large scale seismic velocity heterogeneity in the
27 lowermost mantle [e.g., Grand, 2002], but tomographic studies give the average structure
28 over a depth range of several hundred km. The D'' discontinuity 200-300 km above
29 the core-mantle boundary (CMB) [Lay and Helmberger, 1983], is widely considered to
30 be due to the phase transition between perovskite (pv) and post-perovskite (ppv). As
31 complex structure such as that within the D'' region cannot be resolved by travel-time
32 tomography, it is important to investigate the structure of the lowermost mantle in detail
33 using waveform inversion.

34 We have developed methods to invert seismic waveforms for localized seismic structure
35 and applied these methods to infer the shear wave velocity within D''. We have studied
36 regions with both high [Kawai *et al.*, 2007a, b, 2009] and low [Konishi *et al.*, 2009; Kawai
37 and Geller, 2010] average velocity. The former studies found that high velocities are

38 concentrated in the upper half of the D'' region and that there is a steep negative velocity
39 gradient in the lower half of D''. On the other hand, the latter studies found an "S-shaped"
40 velocity profile, which has a low S-velocity zone in the depth range from 2550 to 2750 km.
41 Both structures can be explained by simple thermal effects for a temperature at the CMB
42 of 3800 K [*Kawai and Tsuchiya, 2009*]. In this study we determine the shear wave velocity
43 structure in the lowermost mantle beneath the Arctic, and compare it to the shear wave
44 velocity immediately above the CMB beneath Central Asia, Central America, and the
45 Pacific to estimate the respective proportion of impurities in Mg-pv or Mg-ppv.

2. Analysis

46 The regions beneath which the structure of D'' can be studied in detail are limited
47 due to the source and receiver geometry. *Kawai et al. [2007b]* studied the shear velocity
48 structure beneath the Arctic using Sd waveforms for many stations for many events as
49 well as a small number of S waveforms. They showed that while long-period Sd data
50 could constrain the average structure in D'' and suggest the existence of two-layered
51 structure, waveforms from stations at epicentral distances $70^\circ < \Delta < 90^\circ$ are required to
52 quantitatively constrain the fine structure within D''.

53 A temporary seismic array, CANOE (CANadian NOrthwest Experiment), was deployed
54 in northwest Canada for two and a half years from 2003 to 2005. The CANOE array
55 provides waveform data at epicentral distances around 80° from intermediate-depth events
56 beneath the Hindu-Kush region. In this study we invert for the fine structure of the
57 lowermost mantle beneath the Arctic using S-wave waveforms, including recordings from

58 the CANOE array. Although CANOE data for Hindu-Kush events have a low S/N ratio,
59 waveform inversion can treat such noisy data.

2.1. Waveform data

60 We use the transverse components of broadband waveform data (obtained by rotating
61 the N-S and E-W components) for 7 events from the IRIS and CNSN data centers (Table
62 1; Fig. 1a). Stations in North America including the CANOE array were chosen (Fig.
63 1b). The CANOE stations recorded waveforms for 3 of the 7 events (shown by asterisks
64 in Table 1). We deconvolve the instrument response and apply a bandpass filter to the
65 data and construct data sets for the period range 8-200 s. We then select records which
66 include data for S, ScS and the other phases which arrive between them at epicentral
67 distances $\Delta < 100^\circ$. We compute the ratio of the maximum amplitude of the data and
68 the corresponding synthetic, and eliminate records for which the ratio is greater than 2
69 or less than 0.5. The dataset consists of 232 records that satisfy the above criteria; 213
70 records which did not satisfy the criteria were rejected. The data are velocity seismograms
71 (with units of m/s after deconvolving the instrument response) with 1 Hz sampling. The
72 reciprocal of the maximum amplitude of each record is used as the weighting factor in the
73 inversion, so that all data have roughly the same importance.

74 The source parameters (moment tensors and centroids) are fixed to the Global CMT
75 solution. *Kawai et al.* [2007a] approximated the source time function as a δ -function at
76 the centroid time for the period range 20-200 s. As this study, however, uses data for
77 the period range 8-200 s, we use boxcar moment rate functions whose half-duration is

78 obtained from the Global CMT solutions. We convolve the boxcar moment rate function
79 with the synthetic seismograms and their partial derivatives in the frequency domain.

80 The waveforms that sample the lowermost mantle pass through the crust, upper mantle,
81 transition zone, and lower mantle, whose effects must be corrected for. We handle this
82 by determining static corrections (time shifts) for the observed waveforms. As only three
83 of the Hindu-Kush events were observed by CANOE stations, this study cannot use the
84 methods of *Kawai and Geller* [2010], whereby the time shifts for events and stations are
85 simultaneously inferred together with the parameters of the Earth model. Following *Fuji*
86 *et al.* [2010], we make static corrections using the time shift that gives the best correlation
87 coefficient between the synthetic and observed seismograms, using the arrivals of direct S
88 waves as the reference.

2.2. Waveform inversion

89 For the details of our inversion methods see *Konishi et al.* [2009], *Kawai and Geller*
90 [2010], *Fuji et al.* [2010], *Fuji* [2010], and the references cited therein. **Our inversion is**
91 **a linearized inversion, with the Earth model in the target zone and time shifts (static**
92 **corrections) to account approximately for the effects of source and structure outside the**
93 **target zone, as the variables. We determine the model parameters by minimizing AIC**
94 **(Akaike's Information Criterion; *Akaike* [1977]), which rewards variance reduction and**
95 **penalizes increases in the complexity of the model.** The initial model is anisotropic PREM
96 [*Dziewonski and Anderson*, 1981]. The source parameters (moment tensors, centroids and
97 half-durations) are fixed to the Global CMT solution. In this study, the S-wave velocities
98 at points above the 'tie-in depth' are fixed to PREM, while those below the tie-in depth are

99 the unknown parameters. We conduct inversions using the first n vectors of the conjugate
100 gradient (CG) decomposition of the matrix of coefficients of the normal equations as the
101 basis functions for the perturbation to the starting model [see *Fuji*, 2010, for details]. We
102 vary two parameters to examine the robustness of the inversion results (Fig. 1c). First,
103 we invert respectively for tie-in depths of 420, 460, and 500 km above the CMB. Second,
104 we invert for models using the first 3 and 4 vectors of the CG decomposition as the basis;
105 these models are labeled CG3 and CG4, respectively.

106 Table 2 shows the variance (computed using the Born approximation) for the various
107 inversions. Defining the variance of the data to be 100%, the variance of the residuals (data
108 – PREM synthetics) is 183.6 %. A further variance reduction (to 127.1 %) is achieved
109 by making the time shift corrections. The variance values for the six models obtained by
110 the CG inversions are in the range from 106.1 % to 107.1 %. Since the velocities at points
111 above the tie-in depth are fixed to the initial model (PREM), the residual error is, at least
112 in part, due to the differences between the real Earth and the Earth model (including Q)
113 outside the study region, and also to source effects. Table 2 also shows the values of AIC
114 for each model. We assume that the effective number of independent data is 12.5 % (1/8)
115 of the above numbers of data points, 16577, because the data are sampled at 1 Hz but
116 are low-passed filtered to exclude periods shorter than 8 s. The lower values of AIC in
117 Table 2 for “PREM with time shift” (as compared to PREM), and for the six CG models
118 (as compared to “PREM with time shift”) show the formal statistical significance of the
119 respective models.

120 We briefly discuss the magnitude of the variance values shown in Table 2. In the case
121 of, for example, a parametric fitting of some curve (e.g., a straight line) to observed
122 data, a variance of over 100 % might appear unacceptably large. But in this case the
123 variance values in Table 2 are for the fit of broadband synthetics and observed waveforms.
124 As shown by the quality control stacks in the on-line supplement (Fig. 2) an acceptable
125 fit is obtained. Furthermore, there is a total of 16,577 data points in the various time
126 series, and only three or four free parameters (the coefficients of the CG basis vectors
127 in the perturbation to the starting model), so (as is confirmed by AIC) the variance
128 reduction for the CG models is statistically significant, notwithstanding the fact that the
129 final variance is on the order of 100 %.

130 Fig. 1c shows that all six of the CG models have basically the same general depth-
131 dependence. There is a velocity increase in the depth range 2500-2700 km and a relative
132 velocity decrease in the depth range from 2700 km to the CMB. The general pattern of
133 the increase and decrease is compatible with our previous results [*Kawai et al.*, 2007b] for
134 D'' beneath the Arctic. As all six models produce roughly the same variance reduction,
135 the differences between these models can be regarded as giving a rough indication of the
136 uncertainty of the inversion results (Fig. 1d). The shaded zone is one nominal standard
137 deviation, taking the six models as independent data. (As is well known, the actual error
138 any well be greater than the nominal error estimates.) The velocity increase and decrease
139 are considerably in excess of the nominal error estimates. *Fuji [2010] conducted a detailed*
140 *analysis of waveform inversion for the upper mantle and mantle transition structure near*
141 *Japan, and demonstrated that despite the large variance for each individual record a*

142 reliable and consistent model was obtained. This conclusion applies to waveform inversion
143 in general.

144 We prepared “quality control stacks,” shown in Fig. 2, for each of the seven events
145 analyzed by this study. These stacks are not intended for use in obtaining the Earth
146 model, but rather as an ancillary check to confirm that the inversion result is reasonable.
147 The stacks are made by aligning and summing the records (after static corrections), using
148 the PREM arrival time and normalizing the maximum amplitude of each of the observed
149 records to one. Because a non-causal filter is used, there are signals before time $t = 0$.
150 The synthetics (computed for model CG4, with a tie-in depth of 500 km) are processed
151 using the same weighting factors as the corresponding observed records.

152 The quality control stacks in Fig. 2 show that the synthetics for the final model are, over-
153 all, a clear improvement over the initial model, thereby confirming that the inversion has
154 reached a reasonable result. The stacked data show ringing (presumably due to secondary
155 scattering due to heterogeneity throughout the Earth) that is not present in the synthetics
156 for either PREM or our final model. Also, there are in some cases significant amplitude
157 differences, presumably due to both anelastic attenuation and focusing/defocusing effects.
158 This figure shows the ability of waveform inversion to obtain robust Earth models from
159 data which are too noisy to permit analysis by trial and error forward modeling.

160 To study the uncertainty of our models and the resolving power of our inversions, and
161 to present further visual examples of observed and synthetic seismograms, additional
162 material is presented in the on-line supplement.

3. Discussion

163 The average composition in the mantle is thought to contain impurities such as alu-
164 minum and iron [*Ringwood, 1962*]. Arguments from mineral physics show that impurities
165 in Mg-pv decrease the shear wave velocity [*Tsuchiya and Tsuchiya, 2006*]. For example,
166 a 1 mol% increase in the amount of both aluminum and iron causes a 0.303 % velocity
167 decrease in Mg-pv or a 0.367 % velocity decrease in Mg-ppv. In this study we estimated
168 the shear wave velocity as about 7.11 km/s immediately above the CMB. As the tem-
169 perature at the CMB is isothermal, the shear wave velocity at the CMB would be equal
170 everywhere if the mineralogical phase and chemical composition is identical [*Kawai and*
171 *Tsuchiya, 2009*]. Since seismic velocity can be directly inferred by waveform inversion,
172 we can use the inferred velocity immediately above the CMB to distinguish the effects of
173 temperature from those of chemical composition.

174 Our previous study of D'' beneath Central America [*Kawai et al., 2007a*] found an S-
175 wave velocity of about 7.25 km/s immediately above the CMB, which is about 0.14 km/s
176 faster than the velocity obtained in this study. This difference can be interpreted as due
177 to the amount of impurities in Mg-ppv beneath the Arctic being 6 mol% larger than
178 that beneath Central America, on the assumption that the ratio of aluminum and iron
179 is constant in Mg-ppv. As PREM can be approximated as pyrolitic composition with
180 7 mol% impurities [*McDonough and Sun, 1995*], we therefore estimate the amount of
181 impurities in Mg-ppv beneath the Arctic to be 13 mol%. *Konishi et al. [2009]* found a
182 shear velocity of 7.18 km/s immediately above the CMB in the western Pacific, which can
183 be interpreted as suggesting that the amount of impurities in Mg-ppv is 3 mol% larger

184 there than that beneath Central America. *Kawai et al.* [2009] found a shear velocity of
185 7.22 km/s at the CMB beneath Central Asia, suggesting that the amount of impurities
186 in Mg-ppv is 1 mol% larger there than beneath central America.

187 Impurities in D'' would cause a two-phase coexistence region for the pv to ppv phase
188 transition. The existence of a two-phase coexistence region has been studied experimen-
189 tally, for example, in the majorite-perovskite system [*Irifune et al.*, 1996; *Hirose et al.*,
190 1999] and the pv-ppv system [*Tateno et al.*, 2007; *Ohta et al.*, 2008]. In multi-component
191 systems in natural compositions such as pyrolite and MORB (mid-ocean ridge basalt) the
192 width of the two-phase coexistence region due to complexities of mineralogy at mantle
193 pressures [*Hirose et al.*, 1999; *Ohta et al.*, 2008] is narrower than that in the MgSiO₃-
194 Al₂O₃ and the MgSiO₃-FeSiO₃ systems [*Ohtani and Sakai*, 2008]. Recent studies on
195 mineral physics show that the width of the two-phase coexistence region is as narrow in
196 the MgSiO₃-Al₂O₃ system [*Tsuchiya and Tsuchiya*, 2008] as seismically detectable dis-
197 continuities.

198 The model obtained in this study shows a velocity increase in the depth range from
199 2500 to 2700 km. This velocity increase is consistent with the experimental result that
200 the post-perovskite phase transition in pyrolitic mantle occurs between 116 and 122 GPa,
201 corresponding to the depth range from 2550 to 2640 km [*Ohta et al.*, 2008]. While the
202 experiments predict a width of about 90 km, our model shows a width of about 200 km.
203 A resolution test (Fig. S2 of the on-line supplement) suggests that a step function dis-
204 continuity would result in an apparent transition depth of 100 km in the seismic velocity
205 models inferred by our methods for the period band used in this study. The 200 km

206 transition in Fig. 1c is thus consistent with a transition over a 100 km depth range, but
207 this is not conclusive. However, when taken together with the extremely low (7.11 km/s)
208 S-velocity just above the CMB, our data are consistent with the presence of significant
209 impurities in D'' beneath the Arctic. Taken into account the resolution of our dataset
210 shown in Fig. S2, we cannot exclude the possibility of a higher temperature in the depth
211 range from 100 km above the CMB to the CMB produced by a hypothetical blanketing
212 effect. As shown in Fig. 1d, the nominal error of the shear velocity immediately above
213 the CMB obtained by this study is about ± 0.02 km/s. If we assume the uncertainty of
214 the velocities immediately above the CMB obtained by the other studies cited above is
215 of comparable order, the difference between the various regions is well in excess of the
216 nominal uncertainty. Further quantification of the uncertainties of the inferred velocities
217 is an important subject for future work.

218 The velocity decrease in the depth range from 2700 km to the CMB can be interpreted
219 as due to the temperature increase in the thermal boundary layer of mantle convection
220 [Kawai and Tsuchiya, 2009], because if a reverse phase transition from ppv to pv occurs
221 due to a temperature increase (a possibility suggested by *Hernlund et al.* [2005]), the shear
222 velocity at the CMB would be even lower (well below 7.0 km/s).

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Event #	Date (Y/M/D)	Latitude	Longitude	Depth	M_w
1	2000/5/1	37.82°	72.43°	152.0	5.6
2	2001/2/25	36.41°	70.62°	193.4	6.1
3*	2004/4/5	36.52°	70.84°	183.5	6.5
4*	2004/8/10	36.52°	70.60°	205.7	6.0
5*	2005/7/23	36.40°	70.65°	212.9	5.5
6	2005/12/12	36.45°	71.06°	210.2	6.5
7	2007/4/3	36.57°	70.59°	221.4	6.2

Table 1. Earthquakes used in this study. The asterisks indicate the three events observed by the CANOE array.

Model	variance (%)	AIC
PREM	183.6	2857.8
PREM with time shift	127.1	2607.1
CG3 (460 km)	106.7	2465.0
CG4 (460 km)	106.1	2462.3
CG3 (480 km)	107.1	2467.6
CG4 (480 km)	106.3	2463.5
CG3 (500 km)	107.1	2468.3
CG4 (500 km)	106.4	2464.5

Table 2. Variance and AIC for each model. The total number of data points is 16577.

Figure 1. (a) Event-receiver geometry, with great circle ray paths. The ray-paths are shown to indicate the coverage, but note that we do not use ray-theoretical approximations. The portions of the great circles which sample D'' are shown in red. Blue reversed triangles and red stars show the sites of stations used in our study and earthquakes studied, respectively. (b) Detailed map of the stations used. CANOE stations are shown by red triangles and permanent stations by blue triangles. (c) The results of inversions for tie-in depths of 420, 460, and 500 km above the CMB for bases of the first 3 or 4 CG vectors, respectively. (d) Nominal error bars (one standard deviation) estimated by treating the six CG models in Fig. 1c as independent.

Figure 2. “Quality control stacks” for each of the seven events, which were computed as follows. First all of the observed waveforms for each event which met the selection criteria were time shifted using PREM, after making the same static corrections as in the inversion. These waveforms were then filtered in the passband 8 s to 200 s using a four-pole non-causal bandpass filter. The maximum amplitude of each observed record was normalized to one, and the waveforms were then stacked (thick curves). The synthetics for the initial model (PREM, dotted curves) and final model (model CG4, for a tie-in depth of 500 km) (thin solid curves) were stacked using the same weighting factors as for the corresponding observed record.